

(Chiral) Virasoro invariance of the tree-level MHV graviton scattering amplitudes

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ABSTRACT: In this paper we continue our study of the tree level MHV graviton scattering amplitudes from the point of view of celestial holography. In [arXiv:2008.04330](https://arxiv.org/abs/2008.04330) we showed that the celestial OPE of two gravitons in the MHV sector can be written as a linear combination of $\overline{\text{SL}(2, \mathbb{C})}$ current algebra and supertranslation descendants. In this note we show that the OPE is in fact manifestly invariant under the infinite dimensional Virasoro algebra as is expected for a 2-D CFT. This is consistent with the conjecture that the holographic dual in 4-D asymptotically flat space time is a 2-D CFT. Since we get only one copy of the Virasoro algebra we can conclude that the holographic dual theory which computes the MHV amplitudes is a chiral CFT with a host of other infinite dimensional global symmetries including $\overline{\text{SL}(2, \mathbb{C})}$ current algebra, supertranslations and subsubleading soft graviton symmetry. We also discuss some puzzles related to the appearance of the Virasoro symmetry.

KEYWORDS: Conformal and W Symmetry, Scale and Conformal Symmetries, Scattering Amplitudes

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1 Introduction and results

In four dimensional asymptotically flat space time the asymptotic symmetries are infinite dimensional [1–11, 13–21] and so in a holographic description the dual theory must have infinite dimensional global symmetry. It has been conjectured [6–10, 21] that the global symmetry group of the dual theory contains the local conformal or Virasoro symmetry. As a result the dual theory is a two dimensional CFT. In this note we study this conjecture in the context of celestial holography [22–63].

We revisit the global symmetries of the tree level MHV graviton scattering amplitudes calculated in General Relativity. In [13] we studied the celestial MHV amplitudes in great detail and showed that the celestial OPE between two graviton primaries in the MHV sector can be written as a linear combination of $\overline{\text{SL}(2, \mathbb{C})}$ current algebra, supetranslations

and global $\text{SL}(2, \mathbb{C})$ descendants. Moreover, due to the existence of the null states [13],¹

$$\Psi_{\Delta}(z, \bar{z}) = \left[J_{-1}^1 P_{-1,-1} - (\Delta - 1) P_{-2,0} \right] G_{\Delta}^{+}(z, \bar{z}) = 0 \quad (1.1)$$

and

$$\Phi_{\Delta}(z, \bar{z}) = \left[L_{-1} P_{-1,-1} + 2J_{-1}^0 P_{-1,-1} - (\Delta + 1) P_{-2,-1} - \bar{L}_{-1} P_{-2,0} \right] G_{\Delta}^{+}(z, \bar{z}) = 0 \quad (1.2)$$

one can write down differential equations which can be solved to determine the MHV graviton scattering amplitudes. Therefore the MHV amplitudes are completely determined by the $\overline{\text{SL}(2, \mathbb{C})}$ current algebra, supertranslations and the global $\text{SL}(2, \mathbb{C})$ symmetries.

In this paper we point out that the celestial OPE of two gravitons in the MHV sector is actually *invariant under the infinite dimensional Virasoro algebra* whose global part is the $\text{SL}(2, \mathbb{C})$ algebra. This shows that the symmetries of the MHV graviton scattering amplitudes is a *semi-direct product of the Virasoro algebra, $\overline{\text{SL}(2, \mathbb{C})}$ current algebra and supertranslations*.² Since we get only one copy of the Virasoro algebra, the *holographic dual theory* which computes the MHV graviton scattering amplitudes is a *chiral (celestial) CFT*.

For the sake of completeness let us now write down the symmetry algebra.

1.1 Virasoro algebra

$$[L_m, L_n] = (m - n)L_{m+n} + \frac{c}{12}n(n^2 - 1)\delta_{m+n,0} \quad m, n \in \mathbb{Z}. \quad (1.3)$$

Although we have written down the central charge term in the Virasoro algebra, our demonstration of Virasoro invariance of the celestial OPE does not allow us to determine the value of the central charge c . So we leave the determination of c to future work.

1.2 $\overline{\text{SL}(2, \mathbb{C})}$ current algebra

$$[J_m^a, J_n^b] = (a - b)J_{m+n}^{a+b}, \quad a, b = 0, \pm 1, \quad m, n \in \mathbb{Z} \quad (1.4)$$

$$J_0^1 = \bar{L}_1, \quad J_0^0 = \bar{L}_0, \quad J_0^{-1} = \bar{L}_{-1}. \quad (1.5)$$

1.3 Supertranslation

$$[P_{m,n}, P_{m',n'}] = 0, \quad n, n' = 0, -1. \quad (1.6)$$

1.4 Mixed commutators

$$[J_m^1, P_{n,-1}] = P_{m+n,0}, \quad [J_m^0, P_{n,-1}] = \frac{1}{2}P_{m+n,-1}, \quad [J_m^{-1}, P_{n,-1}] = 0 \quad (1.7)$$

$$[J_m^1, P_{n,0}] = 0, \quad [J_m^0, P_{n,0}] = -\frac{1}{2}P_{m+n,0}, \quad [J_m^{-1}, P_{n,0}] = -P_{m+n,-1} \quad (1.8)$$

$$[L_m, J_n^a] = -nJ_{m+n}^a, \quad m \in \mathbb{Z}, \quad n \in \mathbb{Z} \quad (1.9)$$

$$[L_n, P_{a,b}] = \left(\frac{n-1}{2} - a \right) P_{a+n,b}, \quad n \in \mathbb{Z}, \quad a \in \mathbb{Z}, \quad b = 0, -1. \quad (1.10)$$

¹The generators of the $\overline{\text{SL}(2, \mathbb{C})}$ current algebra are denoted by J_n^a where $a = 0, \pm 1$ and $n \in \mathbb{Z}$. Similarly we denote the generators of the supertranslations coming from the positive helicity soft graviton by $P_{n,b}$ where $n \in \mathbb{Z}$ and $b = 0, -1$. For further details on the symmetry algebra and null states we refer the reader to [13].

²See [71] for a potential derivation of this symmetry algebra from the bulk gravity point of view.

The symmetry algebra that we have shown can be further extended by including the subsubleading soft graviton symmetry [14–16] and an infinite number of other soft symmetries [14, 16] which appear when the scaling dimension of a (positive helicity) graviton primary assumes negative integer values.

2 Virasoro invariance of the OPE in the MHV sector

Consider the celestial OPE between two positive helicity outgoing gravitons in the MHV sector. This is given by [13],

$$\begin{aligned}
 & G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)\Big|_{\text{MHV}} \\
 &= \frac{\bar{z}}{z} B(\Delta_1 - 1, \Delta_2 - 1) \left[-P_{-1,-1} + z \left(c_1 J_{-1}^0 P_{-1,-1} + c_2 P_{-2,-1} \right) \right. \\
 &\quad \left. + z^2 \left(c_3 J_{-2}^0 P_{-1,-1} + c_4 P_{-3,-1} + c_5 \left(2L_{-1} P_{-2,-1} - 2\bar{L}_{-1} P_{-3,0} + 2L_{-1} \bar{L}_{-1} P_{-2,0} - L_{-1}^2 P_{-1,-1} \right) \right) \right. \\
 &\quad \left. + \mathcal{O}(z^3) \right] G_{\Delta_1 + \Delta_2 - 1}^+(0) + \dots
 \end{aligned} \tag{2.1}$$

The descendant OPE coefficients denoted above by c_i are given by

$$\begin{aligned}
 c_1 &= \frac{2(\Delta_1 - 1)}{\Delta_1 + \Delta_2 - 2}, & c_2 &= -\Delta_1, & c_3 &= \frac{2(\Delta_1 - 1)(\Delta_2 - 1)}{(\Delta_1 + \Delta_2 - 2)(\Delta_1 + \Delta_2 - 1)}, \\
 c_4 &= -\frac{\Delta_1 \Delta_2}{\Delta_1 + \Delta_2 - 1}, & c_5 &= \frac{\Delta_1(\Delta_1 - 1)}{2(\Delta_1 + \Delta_2 - 2)(\Delta_1 + \Delta_2 - 1)}.
 \end{aligned} \tag{2.2}$$

In (2.1) we have, for simplicity, kept only terms of the $\mathcal{O}(\bar{z}z^n)$ where $n \geq -1$ because under the action of the Virasoro generators these terms do not mix with other terms in the OPE.

We already know [13] that the OPE (2.1) is invariant under the action of the global generator L_1 . In this section we show that both sides of the OPE (2.1) transform in the same way under the action of the first Virasoro generator L_2 . This check is non-trivial because the OPE (2.1) has terms of $\mathcal{O}(\bar{z}z)$.

It is important to note that the $\mathcal{O}(\bar{z}z)$ term in (2.1) does not contain the L_{-2} descendant of the graviton primary and so the Virasoro central charge c does not play any role in the process of checking the Virasoro invariance. Let us now spell out the details.

2.1 L_2 invariance

Let us start with the OPE

$$\begin{aligned}
 & G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)\Big|_{\text{MHV}} \\
 &= \frac{\bar{z}}{z} B(\Delta_1 - 1, \Delta_2 - 1) \left[-P_{-1,-1} + z \left(c_1 J_{-1}^0 P_{-1,-1} + c_2 P_{-2,-1} \right) \right. \\
 &\quad \left. + z^2 \left(c_3 J_{-2}^0 P_{-1,-1} + c_4 P_{-3,-1} + c_5 \left(2L_{-1} P_{-2,-1} - 2\bar{L}_{-1} P_{-3,0} + 2L_{-1} \bar{L}_{-1} P_{-2,0} - L_{-1}^2 P_{-1,-1} \right) \right) \right. \\
 &\quad \left. + \mathcal{O}(z^3) \right] G_{\Delta_1 + \Delta_2 - 1}^+(0) + \dots
 \end{aligned} \tag{2.3}$$

We now show that both sides of (2.3) transform in the same way under the action of the Virasoro generator L_2 . For this we take the gravitons in the OPE to be Virasoro primaries so that they satisfy³

$$[L_2, G_{h, \bar{h}}(z, \bar{z})] = z^2(z\partial_z + 3h)G_{h, \bar{h}}(z, \bar{z}). \quad (2.4)$$

Now we apply L_2 to the l.h.s. of (2.3) and obtain using (2.4)

$$[L_2, G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)] = z^2(z\partial_z + 3h_1)G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0) \sim \mathcal{O}(\bar{z}z) + \dots \quad (2.5)$$

where the dots denote terms of $\mathcal{O}(\bar{z}z^2)$ and higher. The leading $\mathcal{O}(\bar{z}z)$ term in (2.5) is given by,

$$\begin{aligned} [L_2, G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)] \Big|_{\mathcal{O}(\bar{z}z)} &= z\bar{z} B(\Delta_1 - 1, \Delta_2 - 1) (1 - 3h_1)P_{-1, -1}G_{\Delta_1 + \Delta_2 - 1}^+(0) \\ &= -z\bar{z} B(\Delta_1 - 1, \Delta_2 - 1) \frac{(3\Delta_1 + 4)}{2} P_{-1, -1}G_{\Delta_1 + \Delta_2 - 1}^+(0). \end{aligned} \quad (2.6)$$

Now we apply to L_2 to the r.h.s. of the OPE (2.3). Using the following commutation relations

$$\begin{aligned} [L_m, J_n^a] &= -nJ_{m+n}^a, & [L_m, P_{r,s}] &= \left(\frac{m-1}{2} - r\right) P_{r+m,s}, \\ [J_m^a, J_n^b] &= (a-b)J_{m+n}^{a+b}, & [P_{m,n}, P_{r,s}] &= 0 \end{aligned} \quad (2.7)$$

and the values of the coefficients c_3, c_4, c_5 given in (2.2) we find that the action of L_2 on the r.h.s. of (2.3) yields

$$\begin{aligned} z\bar{z} B(\Delta_1 - 1, \Delta_2 - 1) \left[(\Delta_1 + \Delta_2 - 2)c_3 + \frac{7c_4}{2} - 3c_5(\Delta_1 + \Delta_2 - 2) \right] P_{-1, -1}G_{\Delta_1 + \Delta_2 - 1}^+(0) + \dots \\ = -z\bar{z} B(\Delta_1 - 1, \Delta_2 - 1) \frac{(3\Delta_1 + 4)}{2} P_{-1, -1}G_{\Delta_1 + \Delta_2 - 1}^+(0) + \dots \end{aligned} \quad (2.8)$$

where the dots denote terms higher order in z . We can see that the $\mathcal{O}(\bar{z}z)$ terms in (2.6) and (2.8) match precisely as expected from the Virasoro invariance of the OPE (2.1).

In the appendix we perform more checks of the Virasoro invariance of the celestial OPE of ++ and +- helicity gravitons in the MHV sector.

2.2 Reorganising the OPE in terms of Virasoro primaries and descendants

We have shown that up to $\mathcal{O}(\bar{z}z)$ terms the OPE of positive helicity outgoing gravitons in the MHV sector given by (2.1) is invariant under the (holomorphic) Virasoro symmetry. This implies that this OPE can be reorganised according to representations of the Virasoro symmetry algebra as is usually done in 2- d CFTs. In other words, we want to show that every term at a particular order in the OPE is either a Virasoro primary or a Virasoro descendant of another Virasoro primary which appears at a lower order in z . The

³Here (h, \bar{h}) denote the holomorphic and antiholomorphic conformal dimensions given by $h = \frac{\Delta + \sigma}{2}$ and $\bar{h} = \frac{\Delta - \sigma}{2}$ where σ is the helicity of the graviton.

OPE coefficient multiplying a Virasoro primary cannot be obtained by using the Virasoro symmetry alone but the OPE coefficient multiplying a Virasoro descendant is uniquely determined by the Virasoro symmetry and the OPE coefficient of the Virasoro primary from which it descends. For simplicity here we reorganise only the $\mathcal{O}(\bar{z}z^n)$ terms in the OPE.

Now in terms of Virasoro primaries and descendants, we find that the positive helicity OPE (2.1) can be expressed as

$$\begin{aligned}
 & G_{\Delta_1}^+(z, \bar{z}) G_{\Delta_2}^+(0) \Big|_{\text{MHV}} \\
 &= \frac{\bar{z}}{z} B(\Delta_1 - 1, \Delta_2 - 1) \left[C_{12\phi_1} \left(1 + z\beta_{\phi_1}^{(1)} L_{-1} + z^2\beta_{\phi_1}^{(1,1)} L_{-1}^2 + z^2\beta_{\phi_1}^{(2)} L_{-2} \right) \phi_1(0) \right. \\
 &\quad \left. + z C_{12\phi_2} \left(1 + z\beta_{\phi_2}^{(1)} L_{-1} \right) \phi_2(0) + z C_{12\phi_3} \left(1 + z\beta_{\phi_3}^{(1)} L_{-1} \right) \phi_3(0) + z^2 C_{12\phi_4} \phi_4(0) + \dots \right] + \dots
 \end{aligned} \tag{2.9}$$

where $\phi_i(0)$ with $i = 1, 2, 3, 4$ are Virasoro primaries. Thus they satisfy the conditions

$$L_0 \phi_i(0) = h_{\phi_i} \phi_i(0), \quad L_m \phi_i(0) = 0, \quad m \geq 1 \tag{2.10}$$

where h_{ϕ_i} is the scaling dimension of the Virasoro primary $\phi_i(0)$.

The coefficients $C_{12\phi_i}$ in (2.9) are given by

$$C_{12\phi_1} = -1, \quad C_{12\phi_2} = \frac{\Delta_1 - 1}{\Delta_1 + \Delta_2 - 2}, \quad C_{12\phi_3} = \frac{\Delta_1 - \Delta_2}{\Delta_1 + \Delta_2 - 2}, \quad C_{12\phi_4} = -\frac{\Delta_1 \Delta_2}{\Delta_1 + \Delta_2 - 1}. \tag{2.11}$$

The $\beta_{\phi_i}^{(\vec{k})}$'s in (2.9) are completely determined by the Virasoro algebra and are given by

$$\begin{aligned}
 \beta_{\phi_i}^{(1)} &= \frac{h_1 - h_2 + h_{\phi_i}}{2h_{\phi_i}}, \\
 \beta_{\phi_i}^{(1,1)} &= \frac{4h_{\phi_i} (2(h_1 - h_2 + h_{\phi_i})^2 + h_2 - 4h_1 - h_{\phi_i}) + c(h_1 - h_2 + h_{\phi_i})(h_1 - h_2 + h_{\phi_i} + 1)}{4h_{\phi_i} (c + 2ch_{\phi_i} + 2h_{\phi_i}(8h_{\phi_i} - 5))} \\
 \beta_{\phi_i}^{(2)} &= \frac{h_{\phi_i}(2h_1 + 2h_2 + h_{\phi_i} - 1) + h_1 + h_2 - 3(h_1 - h_2)^2}{(c + 2ch_{\phi_i} + 2h_{\phi_i}(8h_{\phi_i} - 5))}
 \end{aligned} \tag{2.12}$$

where h_1, h_2 are the holomorphic weights of the primary operators whose OPE is being considered and h_{ϕ_i} is the holomorphic weight of the Virasoro primary $\phi_i(0)$ which appears in this OPE. In our case

$$h_1 = \frac{\Delta_1 + 2}{2}, \quad h_2 = \frac{\Delta_2 + 2}{2}. \tag{2.13}$$

Finally, the explicit expressions for the Virasoro primaries $\phi_i(0)$ are as follows

$$\begin{aligned}
 \phi_1(0) &= P_{-1,-1} G_{\Delta_1 + \Delta_2 - 1}^+(0), \quad \phi_2(0) = \bar{L}_{-1} P_{-2,0} G_{\Delta_1 + \Delta_2 - 1}^+(0), \\
 \phi_3(0) &= \left(P_{-2,-1} - \frac{2}{(\Delta_1 + \Delta_2 + 2)} L_{-1} P_{-1,-1} \right) G_{\Delta_1 + \Delta_2 - 1}^+(0) \\
 \phi_4(0) &= \left(P_{-3,-1} + a_1 J_{-2}^0 P_{-1,-1} + a_2 L_{-2} P_{-1,-1} + a_3 L_{-1}^2 P_{-1,-1} \right. \\
 &\quad \left. + a_4 L_{-1} P_{-2,-1} + a_5 L_{-1} \bar{L}_{-1} P_{-2,0} + a_6 \bar{L}_{-1} P_{-3,0} \right) G_{\Delta_1 + \Delta_2 - 1}^+(0)
 \end{aligned} \tag{2.14}$$

where the coefficients denoted by a_i and b_i are given by

$$\begin{aligned}
 a_1 &= \frac{2(\Delta_1 - 1)(\Delta_2 - 1)}{C_{12\phi_4}(\Delta_1 + \Delta_2 - 2)(\Delta_1 + \Delta_2 - 1)}, & a_2 &= -\frac{C_{12\phi_1} \beta_{\phi_1}^{(2)}}{C_{12\phi_4}}, \\
 a_3 &= \frac{1}{C_{12\phi_4}} \left(-\frac{\Delta_1(\Delta_1 - 1)}{2(\Delta_1 + \Delta_2 - 2)(\Delta_1 + \Delta_2 - 1)} - C_{12\phi_1} \beta_{\phi_1}^{(1,1)} + \frac{2C_{12\phi_3} \beta_{\phi_3}^{(1)}}{(\Delta_1 + \Delta_2 + 2)} \right), \\
 a_4 &= \frac{1}{C_{12\phi_4}} \left(\frac{\Delta_1(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 - 2)(\Delta_1 + \Delta_2 - 1)} - C_{12\phi_3} \beta_{\phi_3}^{(1)} \right), \\
 a_5 &= \frac{1}{C_{12\phi_4}} \left(\frac{\Delta_1(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 - 2)(\Delta_1 + \Delta_2 - 1)} - C_{12\phi_2} \beta_{\phi_2}^{(1)} \right), \\
 a_6 &= -\frac{\Delta_1(\Delta_1 - 1)}{C_{12\phi_4}(\Delta_1 + \Delta_2 - 2)(\Delta_1 + \Delta_2 - 1)}.
 \end{aligned} \tag{2.15}$$

As we have already mentioned we are not able to determine the value of the Virasoro central charge c in this way. The reason being that the $\mathcal{O}(\bar{z}z)$ term in the OPE (2.1) does not contain the L_{-2} descendant of the graviton primary and so both sides of the OPE (2.1) transform in the same way under L_2 irrespective of the value of the central charge c . This is also evident from the decomposition (2.9) which holds for *arbitrary* value of the central charge c . In particular the states $\phi_i(0)$ are Virasoro primaries for any value of the central charge as can be easily checked. We leave the determination of the value of the central charge to future research.⁴

3 Discussion

Since the MHV sector is manifestly Virasoro invariant, the dual theory which computes the MHV scattering amplitudes has a stress tensor $T(z)$ whose modes are the Virasoro generators L_n . In celestial CFT the stress tensor is given by [6, 9, 10] the shadow of the subleading conformally soft graviton. In particular the holomorphic stress tensor $T(z)$ is the shadow of the *negative helicity* subleading soft graviton. Let us try to explain the appearance of $T(z)$ in the MHV sector along this line.

So consider the NMHV amplitudes of the form $\langle - - - + + \dots + \rangle$. Now if we make one of the negative helicity gravitons subleading (conformally) soft [12, 64–70] then the NMHV amplitude soft factorises via the MHV amplitude $\langle - - + + \dots + \rangle$ and this can explain the appearance of the stress tensor $T(z)$ in the MHV sector. But for various reasons this explanation cannot be complete. For example, the negative helicity subleading soft graviton gives rise to $SL(2, \mathbb{C})$ current algebra and its shadow gives the holomorphic stress tensor $T(z)$. Although we find the stress tensor $T(z)$ in the MHV sector, the $SL(2, \mathbb{C})$ current algebra is not visible in the MHV sector. It is not clear that why these two symmetry algebras arising from the negative helicity subleading soft graviton are not treated on an equal footing. Similarly, we also have positive helicity soft graviton in the MHV sector

⁴Since we are considering tree level scattering amplitudes, presumably the value of the central charge c is zero.

and we do find the $\overline{\text{SL}(2, \mathbb{C})}$ current algebra symmetry [13] in the MHV sector but not the $\overline{\text{Virasoro}}$ symmetry. It may be that not all symmetries are manifest [16] at the level of OPE. We hope to return to these issues in future.

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A Virasoro invariance of positive helicity OPE: further checks

In this section, we provide additional checks of the Virasoro invariance of the positive helicity graviton OPE in the MHV sector. In particular we first compute the $\mathcal{O}(z^2)$ terms in the OPE in subsection A.1. The invariance of the OPE at this order under the action of L_2 is proved in subsection A.2. In subsection A.3, we rewrite the terms of $\mathcal{O}(z^n)$ upto $n = 2$ in the positive helicity OPE according to of Virasoro primaries and descendants.

A.1 $\mathcal{O}(z^2)$ terms in positive helicity OPE

Here we compute the $\mathcal{O}(z^2)$ terms in the OPE between positive helicity graviton primaries in the MHV sector. The operators which can appear at this order must have dimensions $\left(\frac{7+h}{2}, \frac{\bar{h}-1}{2}\right)$, where $(h, \bar{h}) = \left(\frac{\Delta_1+\Delta_2+1}{2}, \frac{\Delta_1+\Delta_2-3}{2}\right)$. We then assume the following ansatz for the $\mathcal{O}(z^2)$ OPE⁵

$$\begin{aligned}
 & G_{\Delta_1}^+(z, \bar{z}) G_{\Delta_2}^+(0) \Big|_{\mathcal{O}(z^2)} \\
 &= z^2 B(\Delta_1 - 1, \Delta_2 - 1) \left[x_1 P_{-4,0} + x_2 J_{-3}^1 P_{-1,-1} + x_3 L_{-1}^2 P_{-2,0} + x_4 L_{-1} P_{-3,0} \right] G_{\Delta_1+\Delta_2-1}^+(0)
 \end{aligned} \tag{A.1}$$

where x_1, x_2, x_3, x_4 are as of yet undetermined coefficients. Note that using the null state relations

$$\left[L_{-1} P_{-2,0} + J_{-2}^1 P_{-1,-1} - (\Delta_1 + \Delta_2) P_{-3,0} \right] G_{\Delta_1+\Delta_2-1}^+(0) = 0 \tag{A.2}$$

$$\left[L_{-1} P_{-1,-1} + 2J_{-1}^0 P_{-1,-1} - (\Delta_1 + \Delta_2) P_{-2,-1} - \bar{L}_{-1} P_{-2,0} \right] G_{\Delta_1+\Delta_2-1}^+(0) = 0 \tag{A.3}$$

⁵The ansatz (A.1) is not the most general one, since it is possible to write down additional descendants with dimensions $\left(\frac{7+h}{2}, \frac{\bar{h}-1}{2}\right)$. As we will see in this section, the coefficients x_1, x_2, x_3, x_4 in (A.1) can be uniquely determined using the underlying symmetry algebras. This justifies why it is sufficient to consider (A.1).

we can express (A.1) purely in terms of $\overline{\text{SL}(2, \mathbb{C})}$ current algebra and supertranslation descendants. But for the purpose of determining the coefficients x_i , it turns out to be more convenient to work with the form of the OPE given in (A.1).

We will now compute the x_i 's by demanding that both sides of the OPE transform in the same way under the action of the underlying symmetries. For this, let us first act with the $\overline{\text{SL}(2, \mathbb{C})}$ current algebra generator J_1^0 on both sides of (A.1).

The action of J_1^0 on a conformal primary $\phi_{h, \bar{h}}(z, \bar{z})$ is given by

$$[J_1^0, \phi_{h, \bar{h}}(z, \bar{z})] = z(\bar{z}\partial_{\bar{z}} + \bar{h})\phi_{h, \bar{h}}(z, \bar{z}). \quad (\text{A.4})$$

Using this we get

$$[J_1^0, G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)] = z(\bar{z}\partial_{\bar{z}} + \bar{h}_1)G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0) \quad (\text{A.5})$$

where $\bar{h}_1 = (\Delta_1 - 2)/2$. Then plugging in the $\mathcal{O}(z)$ OPE into the r.h.s. of (A.5) yields

$$\begin{aligned} & [J_1^0, G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)] \Big|_{\mathcal{O}(z^2)} \\ &= z^2 B(\Delta_1 - 1, \Delta_2 - 1) \bar{h}_1 \left[\Delta_1 P_{-3,0} - \frac{(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 - 2)} J_{-2}^1 P_{-1,-1} \right] G_{\Delta_1 + \Delta_2 - 1}^+(0). \end{aligned} \quad (\text{A.6})$$

Now applying J_1^0 to the r.h.s. of (A.1) and using the commutation relations given in (2.7) we obtain

$$\begin{aligned} & [J_1^0, G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)] \Big|_{\mathcal{O}(z^2)} \\ &= z^2 B(\Delta_1 - 1, \Delta_2 - 1) \left[\left(-\frac{x_1}{2} - 2x_4 + (\Delta_1 + \Delta_2)(\Delta_1 + \Delta_2 - 4)x_3 \right) P_{-3,0} \right. \\ & \quad \left. + \left(-x_2 - x_3(\Delta_1 + \Delta_2 - 4) + \frac{x_4}{2} \right) J_{-2}^1 P_{-1,-1} \right] G_{\Delta_1 + \Delta_2 - 1}^+(0). \end{aligned} \quad (\text{A.7})$$

In arriving at (A.7), we have also used the null state relation (A.2). Now invariance of the OPE under J_1^0 implies that (A.6) and (A.7) should match. As a result we get

$$-\frac{x_1}{2} - 2x_4 + (\Delta_1 + \Delta_2)(\Delta_1 + \Delta_2 - 4)x_3 = \bar{h}_1 \Delta_1 \quad (\text{A.8})$$

and

$$x_2 + x_3(\Delta_1 + \Delta_2 - 4) - \frac{x_4}{2} = \frac{\bar{h}_1(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 - 2)}. \quad (\text{A.9})$$

We need a few more equations to solve for all the x_i 's. So, let us act with the global spacetime translation generator $P_{0,-1}$ on both sides of (A.1). $P_{0,-1}$ acts on a outgoing conformal primary $\phi_{h, \bar{h}}(z, \bar{z})$ as follows

$$[P_{0,-1}, \phi_{h, \bar{h}}(z, \bar{z})] = z\phi_{h+1/2, \bar{h}+1/2}(z, \bar{z}). \quad (\text{A.10})$$

(A.10) then implies

$$[P_{0,-1}, G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)] = zG_{\Delta_1+1}^+(z, \bar{z})G_{\Delta_2}^+(0). \quad (\text{A.11})$$

Inserting the $\mathcal{O}(z)$ OPE in the r.h.s. of (A.11) then gives

$$\begin{aligned}
 & [P_{0,-1}, G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)] \Big|_{\mathcal{O}(z^2)} \\
 &= z^2 B(\Delta_1, \Delta_2 - 1) \left[(\Delta_1 + 1)P_{-3,0} - \frac{\Delta_1}{(\Delta_1 + \Delta_2 - 1)} J_{-2}^1 P_{-1,-1} \right] G_{\Delta_1 + \Delta_2}^+(0).
 \end{aligned} \tag{A.12}$$

Now evaluating the action of $P_{0,-1}$ on the r.h.s. of (A.1) we get

$$\begin{aligned}
 & [P_1^0, G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)] \Big|_{\mathcal{O}(z^2)} \\
 &= z^2 B(\Delta_1 - 1, \Delta_2 - 1) \left[(-x_2 + x_4 + 2x_3(\Delta_1 + \Delta_2 + 1))P_{-3,0} - 2x_3 J_{-2}^1 P_{-1,-1} \right] G_{\Delta_1 + \Delta_2}^+(0)
 \end{aligned} \tag{A.13}$$

where we have once again used the null state (A.2). Now (A.12) and (A.13) should match in order for the $\mathcal{O}(z^2)$ OPE to be invariant under $P_{0,-1}$. Therefore, we get

$$x_4 + 2x_3(\Delta_1 + \Delta_2 + 1) - x_2 = \frac{(\Delta_1 + 1)(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 - 2)} \tag{A.14}$$

and

$$x_3 = \frac{\Delta_1(\Delta_1 - 1)}{2(\Delta_1 + \Delta_2 - 1)(\Delta_1 + \Delta_2 - 2)}. \tag{A.15}$$

Equations (A.9) and (A.14) can be simultaneously solved to give

$$x_2 = -\frac{(\Delta_1 - 1)(\Delta_2 - 1)}{(\Delta_1 + \Delta_2 - 1)(\Delta_1 + \Delta_2 - 2)} \tag{A.16}$$

$$x_4 = -\frac{\Delta_1(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 - 1)(\Delta_1 + \Delta_2 - 2)}. \tag{A.17}$$

Plugging these results in (A.8) then yields

$$x_1 = \frac{\Delta_1 \Delta_2}{\Delta_1 + \Delta_2 - 1}. \tag{A.18}$$

The $\mathcal{O}(z^2)$ OPE (A.1) then takes the form

$$\begin{aligned}
 G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0) \Big|_{\mathcal{O}(z^2)} &= z^2 B(\Delta_1 - 1, \Delta_2 - 1) \left[x_1 P_{-4,0} + x_2 J_{-3}^1 P_{-1,-1} \right. \\
 &\quad \left. + x_3 \left(L_{-1}^2 P_{-2,0} - 2L_{-1} P_{-3,0} \right) \right] G_{\Delta_1 + \Delta_2 - 1}^+(0)
 \end{aligned} \tag{A.19}$$

where

$$\begin{aligned}
 x_1 &= \frac{\Delta_1 \Delta_2}{\Delta_1 + \Delta_2 - 1}, & x_2 &= -\frac{(\Delta_1 - 1)(\Delta_2 - 1)}{(\Delta_1 + \Delta_2 - 1)(\Delta_1 + \Delta_2 - 2)}, \\
 x_3 &= \frac{\Delta_1(\Delta_1 - 1)}{2(\Delta_1 + \Delta_2 - 1)(\Delta_1 + \Delta_2 - 2)}.
 \end{aligned} \tag{A.20}$$

A.2 L_2 invariance of $\mathcal{O}(z^2)$ terms in positive helicity OPE

We will now show that the $\mathcal{O}(z^2)$ OPE computed in the previous section and given by (A.19) is invariant under the action of L_2 . In order to do this we start with the Virasoro transformation property of graviton primaries (2.4)

$$[L_2, G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)] = z^2(z\partial_z + 3h_1)G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0). \quad (\text{A.21})$$

Substituting the $\mathcal{O}(z^0\bar{z}^0)$ OPE into the r.h.s. of (A.21) then yields

$$\begin{aligned} [L_2, G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0)] \Big|_{\mathcal{O}(z^2)} &= z^2 B(\Delta_1 - 1, \Delta_2 - 1) 3h_1 P_{-2,0} G_{\Delta_1 + \Delta_2 - 1}^+(0) \\ &= z^2 B(\Delta_1 - 1, \Delta_2 - 1) \frac{3(\Delta_1 + 2)}{2} P_{-2,0} G_{\Delta_1 + \Delta_2 - 1}^+(0). \end{aligned} \quad (\text{A.22})$$

Now using the commutation relations (2.7) and the values of x_1, x_2, x_3 from (A.20) we find that the action of L_2 on the r.h.s. of (A.19) gives⁶

$$\begin{aligned} z^2 B(\Delta_1 - 1, \Delta_2 - 1) \left[\frac{9x_1}{2} + 3x_2(\Delta_1 + \Delta_2 - 2) + 3x_3(\Delta_1 + \Delta_2 - 2) \right] P_{-2,0} G_{\Delta_1 + \Delta_2 - 1}^+(0) \\ = z^2 B(\Delta_1 - 1, \Delta_2 - 1) \frac{3(\Delta_1 + 2)}{2} P_{-2,0} G_{\Delta_1 + \Delta_2 - 1}^+(0). \end{aligned} \quad (\text{A.23})$$

This matches with (A.22) and hence exhibits the Virasoro invariance of the $\mathcal{O}(z^2)$ terms in the OPE (A.19). Let us again note here that as in the $\mathcal{O}(z\bar{z})$ OPE, the L_{-2} descendant of $G_{\Delta_1 + \Delta_2}^+$ does not appear in the OPE (A.19). Therefore our calculation here is insensitive to the central charge of the Virasoro algebra and Virasoro invariance of (A.19) continues to hold for any value of the central charge.

A.3 Reorganising OPE in terms of Virasoro primaries and descendants

In section 2.2 we showed that the $\mathcal{O}(\bar{z}z^n)$ terms in the positive helicity graviton OPE in the MHV sector can be reorganised according to Virasoro primaries and descendants. Here we provide another example of this reorganisation by considering terms of $\mathcal{O}(z^n)$ in the OPE. So, let us begin with the following set of terms in the OPE,

$$\begin{aligned} G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^+(0) \Big|_{\text{MHV}} \\ = B(\Delta_1 - 1, \Delta_2 - 1) \left[P_{-2,0} + z \left(d_1 P_{-3,0} + d_2 J_{-2}^1 P_{-1,-1} \right) \right. \\ \left. + z^2 \left(x_1 P_{-4,0} + x_2 J_{-3}^1 P_{-1,-1} + x_3 \left(L_{-1}^2 P_{-2,0} - 2L_{-1} P_{-3,0} \right) \right) + \mathcal{O}(z^3) \right] G_{\Delta_1 + \Delta_2 - 1}^+(0) + \dots \end{aligned} \quad (\text{A.24})$$

where d_1, d_2 in (A.25) were obtained using symmetries in [13] and are given by

$$d_1 = \Delta_1, \quad d_2 = -\frac{\Delta_1 - 1}{\Delta_1 + \Delta_2 - 2}. \quad (\text{A.25})$$

⁶In order to derive (A.23) we have also used the relation $(J_{-1}^1 P_{-1,-1} - (\Delta - 1)P_{-2,0}) G_{\Delta}^+ = 0$.

x_1, x_2, x_3 are given in (A.20). Now (A.24) can also be written as

$$\begin{aligned}
 & G_{\Delta_1}^+(z, \bar{z}) G_{\Delta_2}^+(0) \Big|_{\text{MHV}} \\
 &= B(\Delta_1 - 1, \Delta_2 - 1) \left[C_{12\psi_1} \left(1 + z\beta_{\psi_1}^{(1)} L_{-1} + z^2\beta_{\psi_1}^{(1,1)} L_{-1}^2 + z^2\beta_{\psi_1}^{(2)} L_{-2} \right) \psi_1(0) \right. \\
 &\quad \left. + zC_{12\psi_2} \left(1 + z\beta_{\psi_2}^{(1)} L_{-1} \right) \psi_2(0) + z^2C_{12\psi_3} \psi_3(0) + \mathcal{O}(z^3) \right] + \dots
 \end{aligned} \tag{A.26}$$

where $\psi_i(0)$ with $i = 1, 2, 3$ are Virasoro primaries for arbitrary values of the central charge. Their explicit expressions are

$$\begin{aligned}
 \psi_1(0) &= P_{-2,0} G_{\Delta_1 + \Delta_2 - 1}^+(0), \quad \psi_2(0) = \left(P_{-3,0} - \frac{3}{(\Delta_1 + \Delta_2 + 4)} L_{-1} P_{-2,0} \right) G_{\Delta_1 + \Delta_2 - 1}^+(0), \\
 \psi_3(0) &= \left(P_{-4,0} + b_1 J_{-3}^1 P_{-1,-1} + b_2 L_{-1} P_{-3,0} + b_3 L_{-1}^2 P_{-2,0} + b_4 L_{-2} P_{-2,0} \right) G_{\Delta_1 + \Delta_2 - 1}^+(0)
 \end{aligned} \tag{A.27}$$

where the b_i 's in (A.27) are given by

$$\begin{aligned}
 b_1 &= -\frac{(\Delta_1 - 1)(\Delta_2 - 1)}{C_{12\psi_3}(\Delta_1 + \Delta_2 - 1)(\Delta_1 + \Delta_2 - 2)} \\
 b_2 &= \frac{1}{C_{12\psi_3}} \left(-\frac{(\Delta_1 - 1)(\Delta_2 - 1)}{(\Delta_1 + \Delta_2 - 1)(\Delta_1 + \Delta_2 - 2)} - C_{12\psi_2} \beta_{\psi_2}^{(1)} \right) \\
 b_3 &= \frac{1}{C_{12\psi_3}} \left(\frac{(\Delta_1 - 1)(\Delta_2 - 1)}{2(\Delta_1 + \Delta_2 - 1)(\Delta_1 + \Delta_2 - 2)} - C_{12\psi_1} \beta_{\psi_1}^{(1,1)} + \frac{3C_{12\psi_2} \beta_{\psi_2}^{(1)}}{(\Delta_1 + \Delta_2 + 4)} \right) \\
 b_4 &= -\frac{C_{12\psi_1} \beta_{\psi_1}^{(2)}}{C_{12\psi_3}}.
 \end{aligned} \tag{A.28}$$

The coefficients $C_{12\psi_i}$ in (A.26) are given by

$$C_{12\psi_1} = 1, \quad C_{12\psi_2} = \frac{\Delta_2 - \Delta_1}{\Delta_1 + \Delta_2 - 2}, \quad C_{12\psi_3} = \frac{\Delta_1 \Delta_2}{\Delta_1 + \Delta_2 - 1}. \tag{A.29}$$

Note that the $C_{12\psi_i}$'s cannot be determined by using only the Virasoro algebra. The coefficients $\beta_{\psi_i}^{(\vec{k})}$ in (A.26) are uniquely fixed by the Virasoro algebra. The explicit formulae for the $\beta_{\psi_i}^{(\vec{k})}$'s take the same form as in (2.12) with h_{ϕ_i} replaced by h_{ψ_i} , where h_{ψ_i} denotes the holomorphic weight of the Virasoro primary ψ_i appearing in the r.h.s. of (A.26).

B Virasoro invariance of mixed helicity OPE

In this section of the appendix, we consider the celestial OPE between an outgoing positive helicity graviton primary and an outgoing negative helicity graviton primary in the MHV sector. Here we show that this mixed helicity OPE is also Virasoro invariant. In order to furnish a non-trivial check of Virasoro invariance we first compute the $\mathcal{O}(z\bar{z})$ terms in the mixed helicity OPE in subsection B.1. This is followed by an explicit check of the invariance of this OPE under the action of L_2 in subsection B.2. In subsection B.3 it is shown that the mixed helicity OPE can be organised according to Virasoro representations.

B.1 $\mathcal{O}(z\bar{z})$ term in mixed helicity OPE using symmetries

The operators which can appear at $\mathcal{O}(z\bar{z})$ in the mixed helicity OPE should have weights $\left(\frac{3+h}{2}, \frac{\bar{h}+1}{2}\right)$, where $(h, \bar{h}) = \left(\frac{\Delta_1+\Delta_2-3}{2}, \frac{\Delta_1+\Delta_2+1}{2}\right)$. Let us then consider the OPE at this order to take the following form

$$\begin{aligned}
 & G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^-(0)\Big|_{\mathcal{O}(z\bar{z})} \\
 &= z\bar{z} B(\Delta_1 - 1, \Delta_2 + 3) \left[\left(\alpha_1 J_{-2}^0 P_{-1,-1} + \alpha_2 P_{-3,-1} + \alpha_3 J_{-1}^0 P_{-2,-1} + \alpha_4 (J_{-1}^0)^2 P_{-1,-1} \right. \right. \\
 &\quad \left. \left. + \alpha_5 \bar{L}_{-1} P_{-3,0} + \alpha_6 \bar{L}_{-1} J_{-1}^0 P_{-2,0} + \alpha_7 J_{-1}^{-1} P_{-2,0} + \alpha_8 \bar{L}_{-1} J_{-2}^1 P_{-1,-1} \right) G_{\Delta_1+\Delta_2-1}^-(0) \right. \\
 &\quad \left. + \left(\alpha_9 P_{-2,-1}^2 + \alpha_{10} \bar{L}_{-1} P_{-2,-1} P_{-2,0} + \alpha_{11} \bar{L}_{-1}^2 P_{-2,0}^2 \right) G_{\Delta_1+\Delta_2-2}^-(0) \right]
 \end{aligned} \tag{B.1}$$

where $\alpha_i, i = 1, 2, \dots, 11$ are as of yet unknown coefficients. In order to determine these using symmetries we follow the same procedure as in section A.1. We impose that both sides of the OPE (B.1) transform identically under the action of supertranslations and the $\overline{\text{SL}}(2, \mathbb{C})$ current algebra. This yields a system of equations for the α_i 's which can then be simultaneously solved. In particular we will apply below the generators $P_{0,-1}, P_{-1,0}$ and J_0^1 for deriving these equations.

First let us consider acting with $P_{0,-1}$. Demanding both sides of (B.1) to transform in the same fashion under the action of $P_{0,-1}$ leads to

$$\alpha_4 = -\frac{2\Delta_1(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)} \tag{B.2}$$

$$-\frac{\alpha_1}{2} - \frac{\alpha_3}{2} + \frac{\alpha_4}{4} = -\frac{(\Delta_1 + 1)(\Delta_1 - 1)}{\Delta_1 + \Delta_2 + 2} \tag{B.3}$$

$$\alpha_6 + 2\alpha_8 = 0. \tag{B.4}$$

Next we act with $P_{-1,0}$ on both sides of the OPE (B.1). In this case we obtain

$$\alpha_6 = -\alpha_4 = \frac{2\Delta_1(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)} \tag{B.5}$$

$$\alpha_8 = -\frac{\Delta_1(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)} \tag{B.6}$$

$$\alpha_{11} = -\frac{\alpha_6}{4} = -\frac{\Delta_1(\Delta_1 - 1)}{2(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)} \tag{B.7}$$

$$\frac{\alpha_1}{2} + \frac{\alpha_4}{4} + \alpha_5 - \alpha_8 = \frac{(\Delta_1 + 1)(\Delta_1 - 1)}{\Delta_1 + \Delta_2 + 2} \tag{B.8}$$

$$\frac{\alpha_3}{2} - \frac{\alpha_6}{2} + \alpha_7 + \alpha_{10} = 0. \tag{B.9}$$

Finally invariance of the OPE under the action of J_0^1 yields

$$\alpha_1 + \alpha_4 + (2\bar{h} - 1)\alpha_8 = -\frac{2\bar{h}_1(\Delta_1 - 1)}{\Delta_1 + \Delta_2 + 2} \tag{B.10}$$

$$\alpha_2 + \alpha_3 + (2\bar{h} - 1)\alpha_5 = 2\bar{h}_1\Delta_1 \tag{B.11}$$

$$\alpha_3 + 2(\Delta + 3)\alpha_4 + (2\bar{h} - 1)\alpha_6 + 2\alpha_7 = 0 \tag{B.12}$$

$$(\Delta + 2)\alpha_3 + 2\alpha_9 + (2\bar{h} - 1)\alpha_{10} = 0 \tag{B.13}$$

$$(\Delta + 2)\alpha_6 + \alpha_{10} + 4(\bar{h} - 1)\alpha_{11} = 0. \tag{B.14}$$

The above system of equations can be solved to uniquely determine all the α_i 's. These are as follows

$$\begin{aligned} \alpha_1 &= \frac{(\Delta_1 - 1)(\Delta_1 + 2\Delta_2 + 6)}{(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)}, & \alpha_2 &= -\Delta_1, & \alpha_3 &= \frac{2\Delta_1(\Delta_1 - 1)}{\Delta_1 + \Delta_2 + 2}, \\ \alpha_4 &= -\frac{2\Delta_1(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)}, & \alpha_5 &= \frac{\Delta_1(\Delta_1 - 1)}{\Delta_1 + \Delta_2 + 2}, & \alpha_6 &= \frac{2\Delta_1(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)} \\ \alpha_7 &= \frac{\Delta_1(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)}, & \alpha_8 &= -\frac{\Delta_1(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)} \\ \alpha_9 &= -\frac{\Delta_1(\Delta_1 - 1)}{2}, & \alpha_{10} &= -\frac{\Delta_1(\Delta_1 - 1)}{\Delta_1 + \Delta_2 + 2}, & \alpha_{11} &= -\frac{\Delta_1(\Delta_1 - 1)}{2(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)}. \end{aligned} \tag{B.15}$$

B.2 L_2 invariance of $\mathcal{O}(z\bar{z})$ terms in mixed helicity OPE

We shall now verify that the mixed helicity OPE at $\mathcal{O}(z\bar{z})$ (B.1) is invariant under L_2 . For this we act with L_2 on both sides of (B.1). From the l.h.s. we get

$$[L_2, G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^-(0)] \Big|_{\mathcal{O}(z\bar{z})} = -z\bar{z} B(\Delta_1 - 1, \Delta_2 + 3) \frac{(3\Delta_1 + 4)}{2} P_{-1, -1} G_{\Delta_1 + \Delta_2 - 1}^-(0). \tag{B.16}$$

The action of L_2 on the r.h.s. of (B.1) evaluates to

$$z\bar{z} B(\Delta_1 - 1, \Delta_2 + 3) \left[(\Delta_1 + \Delta_2 + 2)\alpha_1 + \frac{7\alpha_2}{2} + \alpha_5 - \frac{\alpha_6}{2} \right] P_{-1, -1} G_{\Delta_1 + \Delta_2 - 1}^-(0). \tag{B.17}$$

Using the values of $\alpha_1, \alpha_2, \alpha_5, \alpha_6$ from (B.15), we find that (B.17) simplifies to

$$-z\bar{z} B(\Delta_1 - 1, \Delta_2 + 3) \frac{(3\Delta_1 + 4)}{2} P_{-1, -1} G_{\Delta_1 + \Delta_2 - 1}^-(0). \tag{B.18}$$

This matches with (B.16) and thereby proves that the $\mathcal{O}(z\bar{z})$ term in the mixed helicity OPE is Virasoro invariant. As in the case of the positive helicity OPE, the central charge of the Virasoro algebra does not play any role here.

B.3 Reorganising OPE in terms of Virasoro primaries and descendants

In the MHV sector, the mixed helicity graviton OPE takes the following form

$$\begin{aligned} & G_{\Delta_1}^+(z, \bar{z})G_{\Delta_2}^-(0) \Big|_{\text{MHV}} \\ &= \frac{\bar{z}}{z} B(\Delta_1 - 1, \Delta_2 + 3) \left[-P_{-1, -1} G_{\Delta_1 + \Delta_2 - 1}^-(0) + z(\rho_1 J_{-1}^0 P_{-1, -1} + \rho_2 P_{-2, -1}) G_{\Delta_1 + \Delta_2 - 1}^-(0) \right. \\ & \quad \left. + z^2(\alpha_1 J_{-2}^0 P_{-1, -1} + \alpha_2 P_{-3, -1} + \alpha_5(2J_{-1}^0 P_{-2, -1} + \bar{L}_{-1} P_{-3, 0})) \right] \end{aligned}$$

$$\begin{aligned}
 & + \frac{\alpha_6}{2} \left(-2(J_{-1}^0)^2 P_{-1,-1} + 2\bar{L}_{-1} J_{-1}^0 P_{-2,0} + J_{-1}^{-1} P_{-2,0} - \bar{L}_{-1} J_{-2}^1 P_{-1,-1} \right) G_{\Delta_1+\Delta_2-1}^-(0) \\
 & + z^2 \left(\alpha_9 P_{-2,-1}^2 + \alpha_{10} \bar{L}_{-1} P_{-2,-1} P_{-2,0} + \alpha_{11} \bar{L}_{-1}^2 P_{-2,0}^2 \right) G_{\Delta_1+\Delta_2-2}^-(0) + \mathcal{O}(z^3) \Big] \\
 & + B(\Delta_1 - 1, \Delta_2 + 3) \left[P_{-2,0} G_{\Delta}^-(0) + z \left(\eta_1 P_{-3,0} + \eta_2 J_{-2}^1 P_{-1,-1} \right) G_{\Delta_1+\Delta_2-1}^-(0) + \mathcal{O}(z^2) \right] + \dots
 \end{aligned} \tag{B.19}$$

where ρ_1, ρ_2 and η_1, η_2 are given by

$$\rho_1 = \frac{2(\Delta_1 - 1)}{\Delta_1 + \Delta_2 + 2}, \quad \rho_2 = -\Delta_1, \quad \eta_1 = \Delta_1, \quad \eta_2 = -\frac{\Delta_1 - 1}{\Delta_1 + \Delta_2 + 2}. \tag{B.20}$$

The values of ρ_1, ρ_2 were computed previously in [13]. Although we have not shown in this paper, η_1, η_2 can also be easily obtained from symmetries. The α_i 's are given by (B.15). The dots in (B.19) denote further terms which we have suppressed for simplicity.

Now it is again possible to rearrange (B.19) in terms of Virasoro primaries and descendants, so as to make the Virasoro invariance of the OPE manifest. Carrying out this reorganisation we find the following result

$$\begin{aligned}
 & G_{\Delta_1}^+(z, \bar{z}) G_{\Delta_2}^-(0) \Big|_{\text{MHV}} \\
 & = \frac{\bar{z}}{z} B(\Delta_1 - 1, \Delta_2 + 3) \left[C_{12\varphi_1} \left(1 + z\beta_{\varphi_1}^{(1)} L_{-1} + z^2\beta_{\varphi_1}^{(1,1)} L_{-1}^2 + z^2\beta_{\varphi_1}^{(2)} L_{-2} \right) \varphi_1(0) \right. \\
 & \quad \left. + z C_{12\varphi_2} \left(1 + z\beta_{\varphi_2}^{(1)} L_{-1} \right) \varphi_2(0) + z C_{12\varphi_3} \left(1 + z\beta_{\varphi_3}^{(1)} L_{-1} \right) \varphi_3(0) + z^2 C_{12\varphi_4} \varphi_4(0) + \mathcal{O}(z^3) \right] \\
 & \quad + B(\Delta_1 - 1, \Delta_2 + 3) \left[C_{12\Omega_1} \left(1 + z\beta_{\Omega_1}^{(1)} L_{-1} \right) \Omega_1(0) + z C_{12\Omega_2} \Omega_2(0) + \mathcal{O}(z^2) \right] + \dots
 \end{aligned} \tag{B.21}$$

where the operators denoted by $\varphi_i(0)$ and $\Omega_i(0)$ are Virasoro primaries (for any value of the central charge). Their explicit expressions are given by (B.23) and (B.24). The coefficients $C_{12\varphi_i}$ and $C_{12\Omega_i}$ are

$$C_{12\varphi_1} = -1, \quad C_{12\varphi_2} = -\Delta_1, \quad C_{12\varphi_3} = -\Delta_1, \quad C_{12\Omega_1} = 1, \quad C_{12\Omega_2} = \Delta_1. \tag{B.22}$$

Once again it is worth noting that Virasoro symmetry alone does not fix the values given in (B.22). The coefficients $\beta_{\varphi_i}^{(\vec{k})}$'s and $\beta_{\Omega_i}^{(\vec{k})}$'s in (B.21), which are uniquely determined by the Virasoro algebra, are given by similar expressions as in (2.12). Let us finally note the expressions for the Virasoro primaries $\varphi_i(0), \Omega_i(0)$.

$$\begin{aligned}
 \varphi_1(0) & = P_{-1,-1} G_{\Delta_1+\Delta_2-1}^-(0), \\
 \varphi_2(0) & = \left(P_{-2,-1} - \frac{2(\Delta_1 - 1)}{\Delta_1(\Delta_1 + \Delta_2 + 2)} J_{-1}^0 P_{-1,-1} - \frac{(\Delta_1 + 1)}{\Delta_1(\Delta_1 + \Delta_2 - 2)} L_{-1} P_{-1,-1} \right) G_{\Delta_1+\Delta_2-1}^-(0) \\
 \varphi_3(0) & = \left[P_{-3,-1} - \frac{(\Delta_1 - 1)(\Delta_1 + 2\Delta_2 + 6)}{\Delta_1(\Delta_1 + \Delta_2 + 3)} J_{-2}^0 P_{-1,-1} + \right. \\
 & \quad \left. - \frac{(\Delta_1 - 1)}{\Delta_1 + \Delta_2 + 2} (2J_{-1}^0 P_{-2,-1} + \bar{L}_{-1} P_{-3,0}) \right]
 \end{aligned}$$

$$\begin{aligned}
& + \frac{(\Delta_1 - 1)}{(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)} (2(J_{-1}^0)^2 P_{-1,-1} - 2\bar{L}_{-1} J_{-1}^0 P_{-2,0} + \\
& - J_{-1}^{-1} P_{-2,0} + \bar{L}_{-1} J_{-2}^1 P_{-1,-1}) \Big] G_{\Delta_1 + \Delta_2 - 1}^-(0) \\
& + \left(\frac{(\Delta_1 - 1)}{2} P_{-2,-1}^2 + \frac{(\Delta_1 - 1)}{\Delta_1 + \Delta_2 + 2} \bar{L}_{-1} P_{-2,-1} P_{-2,0} \right. \\
& + \left. \frac{(\Delta_1 - 1)}{2(\Delta_1 + \Delta_2 + 2)(\Delta_1 + \Delta_2 + 3)} \bar{L}_{-1}^2 P_{-2,0}^2 \right) G_{\Delta_1 + \Delta_2 - 2}^-(0) + \\
& - \frac{(\Delta_1 - 4)(4\Delta_1 + 3)}{2\Delta_1(\Delta_1 + \Delta_2 - 2)(4\Delta_1 + \Delta_2 - 13)} L_{-1}^2 P_{-1,-1} G_{\Delta_1 + \Delta_2 - 1}^-(0) + \\
& - \frac{(\Delta_1(3\Delta_2 - 8) + 4\Delta_2 - 10)}{\Delta_1(\Delta_1 + \Delta_2 - 2)(4\Delta_1 + \Delta_2 - 13)} L_{-2} P_{-1,-1} G_{\Delta_1 + \Delta_2 - 1}^-(0) - \frac{(\Delta_1 + 2)}{\Delta_1 + \Delta_2} L_{-1} \varphi_2(0)
\end{aligned} \tag{B.23}$$

and

$$\begin{aligned}
\Omega_1(0) &= P_{-2,0} G_{\Delta_1 + \Delta_2 - 1}^-(0), \\
\Omega_2(0) &= \left(P_{-3,0} - \frac{(\Delta_1 + 2)}{\Delta_1(\Delta_1 + \Delta_2)} L_{-1} P_{-2,0} - \frac{(\Delta_1 - 1)}{\Delta_1(\Delta_1 + \Delta_2 + 2)} J_{-2}^1 P_{-1,-1} \right) G_{\Delta_1 + \Delta_2 - 1}^-(0).
\end{aligned} \tag{B.24}$$

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